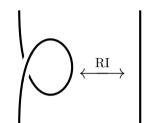
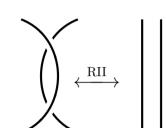
The Jones polynomial in terms of braid group representations

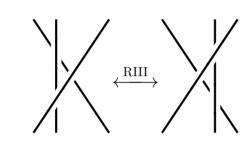
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Overview

A knot is a simple closed curve in \mathbb{R}^3 ; we draw a knot diagram by projecting a knot onto \mathbb{R}^2 . A link is a disjoint union of knots, each considered a component of the link. Two knots are isotopic if one can be continuously deformed into each other without 'breaking' the knot, or equivalently, if they are related to another by a sequence of Reidemeister moves (RI, RII, RIII):







A knot invariant assigns a value to a knot such that if two knots are isotopic, they have the same value under the invariant. We study the Jones polynomial knot invariant and outline a construction of the Jones polynomial via representations of braid groups.

The Kauffman bracket and Jones polynomial

For a link diagram D, the Kauffman bracket $\langle D \rangle$ is a Laurent polynomial in A defined recursively by the relations

$$\left\langle \sum_{i} \right\rangle = A \left\langle \right\rangle \left(\right) + A^{-1} \left\langle \sum_{i} \right\rangle, \quad \left\langle \bigcap_{i} D \right\rangle = (-A^2 - A^{-2}) \left\langle D \right\rangle$$

$$\left\langle \varnothing \right\rangle = 1$$

where the diagrams in the first relation correspond to links identical everywhere except around a crossing, and the diagram of the second relation corresponds to the nonintersecting union of an unknot and D.

By orienting each component of a link D, we define the sign of a crossing and the writhe w(D) of an oriented link D as follows:

$$\sum$$
 = positive crossing, \sum = negative crossing,

$$w(D) = (\#positive crossings) - (\#negative crossings).$$

Then by modifying the Kauffman bracket to be invariant under the Reidemeister moves, we obtain the Kauffman bracket invariant $(-A^3)^{-w(D)}\langle D\rangle$ for an oriented link. The Jones polynomial $V_L(t)$ can be defined in terms of the Kauffman bracket by dividing the value of the Kauffman bracket invariant by $(-A^2 - A^{-2})$ and substituting $A^2 = t^{-1/2}$.

The skein relation of the Jones polynomial is

$$t^{-1}V_{L_{+}}(t) - tV_{L_{-}}(t) = (t^{1/2} - t^{-1/2})V_{L_{0}}(t)$$

where L_+ , L_- and L_0 are link diagrams identical everywhere except for an area around a crossing, replaced by a positive crossing, negative crossing and no crossing respectively.

The braid group and Markov moves

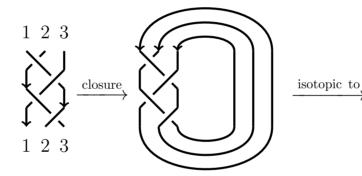
A braid is a set of n strands in $\mathbb{R}^2 \times \{0,1\}$, each with distinct endpoints: one in $\{1,2,...,n\} \times \{0\} \times \{0,1\}$ and one in $\{1,2,...,n\} \times \{1\} \times \{0,1\}$. Any braid can be expressed as an element of the braid group of n strands B_n , with presentation

generators:
$$\sigma_1, \sigma_2, ..., \sigma_{n-1}$$

relations:
$$\sigma_i \sigma_j = \sigma_j \sigma_i, |i - j| \ge 2,$$

$$\sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1},$$

where σ_i denotes a positive crossing between the *i*-th and (i+1)-th strands. The closure of a braid can be formed by connecting the ends (i,1,0) and (i,1,1); by Alexander's theorem, any (oriented) link is isotopic to the closure of a braid. For example, the closure of the braid $\sigma_1^{-1}\sigma_2\sigma_1^{-1}\sigma_2$ is isotopic to the figure-eight knot as shown:



The closures of braids a and b are isotopic to each other if and only if braids a and b are related to each other by the Markov moves (MI, MII):

$$MI: ab \longleftrightarrow ba, MII: b \longleftrightarrow b\sigma_n^{\pm 1},$$

where $a, b \in B_n$ and $b\sigma_n^{\pm 1} \in B_{n+1}$.

Representations of the braid group

For a vector space V over \mathbb{C} , we consider the representation ψ_n : $B_n \to \operatorname{End}(V^{\otimes n})$ given by $\psi_n(\sigma_i) = (id_V)^{\otimes (i-1)} \otimes R \otimes (id_V)^{\otimes (n-i-1)}$, where $R \in \operatorname{End}(V \otimes V)$. From the presentation of the braid group, ψ_n must satisfy

$$\psi_n(\sigma_i)\psi_n(\sigma_j) = \psi_n(\sigma_j)\psi_n(\sigma_i), |i-j| \ge 2,$$

$$\psi_n(\sigma_i)\psi_n(\sigma_{i+1})\psi_n(\sigma_i) = \psi_n(\sigma_{i+1})\psi_n(\sigma_i)\psi_n(\sigma_{i+1}).$$

From these relations, R must satisfy the Yang-Baxter equation

$$(id_V \otimes R)(R \otimes id_V)(id_V \otimes R) = (R \otimes id_V)(id_V \otimes R)(R \otimes id_V).$$

Representations of the braid group (cont.)

We consider the case where V is a 2-dimensional vector space and R is of the form

$$R = \begin{bmatrix} a & 0 & 0 & 0 \\ 0 & b & c & 0 \\ 0 & d & e & 0 \\ 0 & 0 & 0 & f \end{bmatrix}.$$

From the Yang-Baxter equation, if b = 0 and $e \neq 0$, after normalising the entries in terms of t we obtain the R matrices

$$R = egin{bmatrix} t^{1/2} & 0 & 0 & 0 \ 0 & 0 & t & 0 \ 0 & t & t^{1/2} - t^{3/2} & 0 \ 0 & 0 & 0 & t^{1/2} \end{bmatrix}, egin{bmatrix} t^{-1/2} & 0 & 0 & 0 \ 0 & 0 & 1 & 0 \ 0 & 1 & t^{-1/2} - t^{1/2} & 0 \ 0 & 0 & 0 & -t^{1/2} \end{bmatrix}.$$

which correspond to the Jones and Alexander polynomials respectively. By taking the trace of $\psi_n(b)$ where b is a braid, we try to construct an isotopy invariant of a link, which requires $\operatorname{tr}(\psi_n(b))$ to be invariant under the Markov moves:

$$\operatorname{tr}(\psi_n(ab)) = \operatorname{tr}(\psi_n(ba)),$$
$$\operatorname{tr}(\psi_n(b)) = \operatorname{tr}(\psi_{n+1}(b\sigma_n^{\pm 1}))$$

However, this fails, since we require $\operatorname{tr}_2 R^{\pm 1} = i d_V$ but $\operatorname{tr}_2 R^{-1} \neq i d_V$. By using $h = \begin{bmatrix} t^{-1/2} & 0 \\ 0 & t^{1/2} \end{bmatrix}$, and using the properties $R \cdot (h \otimes h) = (h \otimes h) \cdot R$ and $\operatorname{tr}_2((i d_V \otimes h) \cdot R^{\pm 1}) = i d_V$ obtained by construction of R, we obtain the invariant $\operatorname{tr}(h^{\otimes n} \cdot \psi_n(b))$, since

$$\operatorname{tr}(h^{\otimes n} \cdot \psi_n(ab)) = \operatorname{tr}(h^{\otimes n} \cdot \psi_n(a)\psi_n(b)) = \operatorname{tr}(\psi_n(b) \cdot h^{\otimes n} \cdot \psi_n(a))$$
$$= \operatorname{tr}(h^{\otimes n} \cdot \psi_n(b)\psi_n(a)) = \operatorname{tr}(h^{\otimes n} \cdot \psi_n(ba)) \text{ and }$$

$$\operatorname{tr}(h^{\otimes (n+1)} \cdot \psi_{n+1}(\sigma_n^{\pm 1}b)) = \operatorname{tr}(h^{\otimes (n+1)} \cdot (id_V^{\otimes (n-1)} \otimes R^{\pm 1}) \cdot \psi_{n+1}(b))$$
$$= \operatorname{tr}(h^{\otimes n} \cdot \psi_n(b)).$$

The skein relation of this invariant is the same as that of the Jones polynomial since

$$t^{-1}R - tR^{-1} = (t^{-1/2} - t^{1/2})id_{V \otimes V}$$

and R, R^{-1} and $id_{V\otimes V}$ correspond to the three types of crossings. Since the value of the unknot under this invariant is $t^{1/2} + t^{-1/2}$, its value is $t^{1/2} + t^{-1/2}$ times the Jones polynomial $V_L(t)$. This gives an alternate definition of the Jones polynomial in terms of a braid b and its closure L:

$$\operatorname{tr}(h^{\otimes n} \cdot \psi_n(b)) = (t^{1/2} + t^{-1/2})V_L(t)$$

References

Ohtsuki, T. (2002). Quantum Invariants: A Study of Knots, 3-Manifolds, and their Sets. World Scientific. Jones, V.F.R. (1987). Hecke algebra representations of braid groups and link polynomials. Annals of Mathematics. 126 (2):335-388.

Alexander, J. (1923). A lemma on a system of knotted curves. Proceedings of the National Academy of Sciences of the United States of America. 9 (3):93-95.

Murasugi, K. (2008). Knot Theory and its Applications. Reprint of the 1996 edition. Birkhäuser.