

## Introduction

Quantum groups were first introduced in the 1980s by Drinfeld and Jimbo, as a way to solve the Yang-Baxter equation. They are noncommutative algebras that can be deformed. They were independently constructed from a mathematical physics perspective, and one of statistical physics. This project explored three different quantum groups, and worked to identify their application in quantum mechanics.

FIGURE 1: AN ARTISTS IMPRESSION OF QUANTUM PHYSICS



## Methodology:

$$U_q \rightarrow R \rightarrow M \rightarrow T \rightarrow H$$

FIGURE 2: DIAGRAM DEPICTING THE LINK BETWEEN QUANTUM GROUPS AND QUANTUM MECHANICS

The universal R matrix is a solution to the Yang-Baxter equation derived from a quantum group. A monodromy matrix is derived from the R matrix, then the transfer matrix. Finally, the Hamiltonian for a quantum system is formed.

$U_q$

The generators and commutation relations for each group are shown below:

- The Heisenberg algebra:

$$\{a_n, a_{-n}\}_{n>0} \\ [a_n, a_{-m}] = \delta_{n,m} n$$

- Quantum  $sl_2$ :

$$e, f, k = q^h \\ kek^{-1} = q^2 e, \quad kfk^{-1} = q^{-2} f, \quad [e, f] = \frac{k - k^{-1}}{q - q^{-1}}$$

- Quantum affine  $sl_2$ :

$$e_0, f_0, e_1, f_1, k_0, k_1 \\ q^x e_i q^{-x} = q^{\alpha_i(x)} e_i, \quad q^x f_i q^{-x} = q^{-\alpha_i(x)} f_i \\ [e_i, f_j] = \frac{1}{q - q^{-1}} \delta_{ij} (q^{h_i} - q^{-h_i})$$

$R$

The calculation of the R matrix is through the use of a Drinfeld double, where  $a_i$  are basis elements.

$$R = \sum_i a_i \otimes a_i^* \quad \begin{matrix} R \in A \hat{\otimes} A^\circ \\ a_i \in A \\ a_i^* \in A^\circ \end{matrix}$$

The matrix obtained in this fashion obeys the following relations:

$$R\Delta(a) = \Delta^{\circ p}(a)R \\ (1 \otimes \Delta)R = R_{13}R_{23}, \quad (\Delta \otimes 1)R = R_{13}R_{23}$$

These properties alone are enough to prove that R is a solution to the Yang-Baxter Equation.

$$R_{12}R_{13}R_{23} = R_{23}R_{13}R_{12}$$

$M, T \text{ \& } H$

The following relations depict how one can obtain the Hamiltonian from the universal R matrix for periodic boundary conditions. When  $T$  has been diagonalised, so has  $H$ . This is incredibly useful when it is difficult to diagonalise  $H$  directly.

$$M_N(w) = \Pi_{k=1}^N R_{0k}\left(\frac{z_k}{w}\right) \\ T(w) = \text{Tr}_0(M_N(w)) \\ H := T^{-1}(1) \left. \frac{\partial}{\partial w} T(w) \right|_{w=1}$$

## Results:

$$R = \begin{bmatrix} qz - q^{-1}z^{-1} & 0 & 0 & 0 \\ 0 & z - z^{-1} & (q - q^{-1})z & 0 \\ 0 & (q - q^{-1})z^{-1} & z - z^{-1} & 0 \\ 0 & 0 & 0 & qz - q^{-1}z^{-1} \end{bmatrix}$$

FIGURE 3:  $R$ , A REPRESENTATION OF THE UNIVERSAL R MATRIX OF  $U_q(\mathfrak{sl}_2)$  WITH  $z$  THE SPECTRAL PARAMETER.

$$H = -\frac{2}{q - q^{-1}} \begin{bmatrix} q + q^{-1} + 1 & 0 & 0 & 0 \\ 0 & 2 & 1 & 0 \\ 0 & 1 & 2q & 0 \\ 0 & 0 & 0 & q + q^{-1} + 1 \end{bmatrix}$$

FIGURE 4:  $H$ , A HEISENBERG MATRIX CORRESPONDING TO  $M_2$

## Discussion:

The final Hamiltonian matrix found in figure 4 is a great example of how a method like this can work. The use of these theoretical algebra structures to learn information about physical quantum systems is incredibly fascinating and useful. The three groups studied in the scope of this project each had their unique properties, yet there are many more quantum groups out there with links to other physical systems. A more comprehensive review would explore these. There is always room for improvement: figure 4 has yet to be diagonalised. Hence, further research into techniques of diagonalisation, and computational methods to find the Heisenberg matrix using the monodromy matrix with a general N would be informative.

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